

# On the use of the Adjoint Operator in Estimating Neutral Particles Sources

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# Motivations

The **adjoint flux**, the solution of the **adjoint transport equation**, is an important **mathematical tool**. In neutral particles transport theory, some applications may benefit from its study, such as

- ▶ Nuclear Reactors Analysis [4]:
  - ▶ Importance function;
  - ▶ Criticality problems.
- ▶ **Inverse Problems** [5, 6]:
  - ▶ Applications such as: nuclear safety, non-destructive material identification;
  - ▶ In such problems, we are often concerned with **estimating the spatial distribution and intensity of internal sources** of neutral particles;
  - ▶ The **adjoint flux** frequently appear in iterative schemes, when computing gradients related to minimization procedures, or in the evaluation of detector readings in the so-called **source-detector problems**.

This way, we need to obtain accurate solutions to the adjoint of transport equation.

- ▶ The **ADO method** (*Analytical Discrete Ordinates* method) was developed and successfully applied to solve problems in:
  - ▶ radiative heat transfer [3];
  - ▶ rarefied gases dynamic [9];
  - ▶ neutral particles transport [1].
- ▶ The methodology was extended to the **one-dimensional** planar geometry **adjoint** of the transport equation, considering arbitrarily high anisotropy degree [8];
- ▶ Then, the formulation was further extended to **two-dimensional** XY-geometry transport problems in a medium with isotropic scattering [2];
- ▶ In such problems, we were able to derive **closed form expressions** to the averaged fluxes in either  $x$  or  $y$  directions.

- ▶ In this work, the adjoint flux is applied in **estimating internal sources** of particles, **from measurements of internal particles detectors**:
  - ▶ We use the **iterated Tikhonov** (IT) algorithm [6];
  - ▶ And the **Markov Chain Monte Carlo** (MCMC) method within the Bayesian framework.
- ▶ We consider that all **physical properties** of the medium and **incoming fluxes** at the boundaries **are known**;
- ▶ Then, we derive using the adjoint of the transport equation a **linear model** that relates **absorption rate measurements** of a series of particle detectors with the coefficients of the **source expansion in a given basis**.

# Transport Equation

# Transport Equation

- ▶ In steady state, considering **two-dimensional** Cartesian geometry ( $XY$ -geometry), considering **isotropic scattering** and the **monoenergetic approximation**, the transport equation takes the form:

$$\boldsymbol{\Omega} \cdot \nabla \psi + \sigma \psi = \int_{\mathbb{S}'} \sigma_s \psi \, d\boldsymbol{\Omega}' + S,$$

with  $(x, y) \in D = [0, X] \times [0, Y]$  and  $\psi = \psi(x, y, \eta, \xi)$ ;

- ▶ This way, we can rewrite the **transport equation** in a more compact form as:

$$\mathcal{L}\psi = S.$$

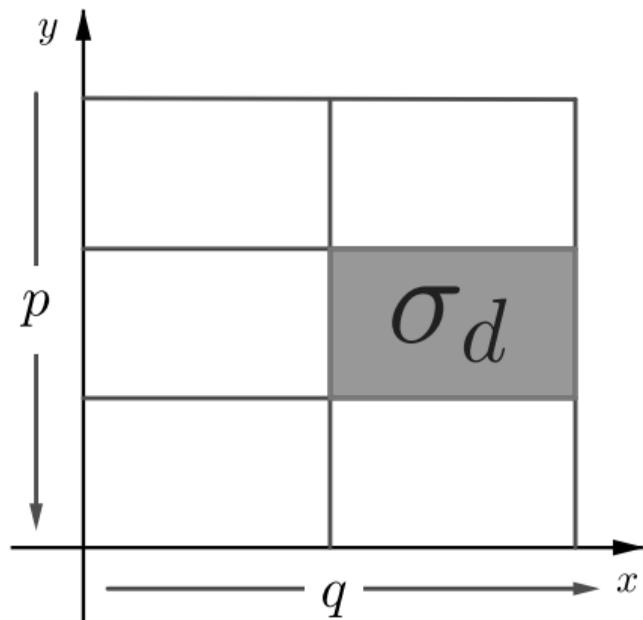
- ▶ We also prescribe **boundary conditions** at the **inward** directions of the domain:

$$\psi(\cdot, \boldsymbol{\Omega}_{in}) = \psi_b(\cdot, \boldsymbol{\Omega}_{in}) + \rho\psi(\cdot, \boldsymbol{\Omega}_{out});$$

# Absorption Rate Measurements

## Absorption Rate Measurements

- ▶ Rectangular mesh of  $n_x$  by  $n_y$  rectangles;
- ▶ Internal **particles detectors** with absorption cross section  $\sigma_d$ :



$$\sigma_d(x, y) = \begin{cases} \sigma_{d,p,q}, & (x, y) \in [x_q, x_{q+1}] \times [y_{p+1}, y_p] \\ 0, & \text{otherwise} \end{cases}$$

# Absorption Rate Measurements

- ▶ We may **evaluate** the absorption rate of particles migrating from all directions within a region  $R$  of the domain using the expression:

$$r = \langle \sigma_d, \psi \rangle;$$

- ▶ Where  $\langle \cdot, \cdot \rangle$  is the **inner product** in the phase space:

$$\langle f, g \rangle = \int_{\mathbb{S}^1} \int_V f(\mathbf{r}, \boldsymbol{\Omega}) g(\mathbf{r}, \boldsymbol{\Omega}) dV d\boldsymbol{\Omega};$$

- ▶ **Remember:** Our goal is to estimate the **spatial distribution and magnitude** of an isotropic internal source of particles  $S$  from a series of known measurements provided by  $n_d$  **particles detectors**, that is,

$$\mathbf{r} = [ \overbrace{r_1}^{\text{detector 1}} \quad \overbrace{r_2}^{\text{detector 2}} \quad \dots \quad \overbrace{r_{n_d}}^{\text{detector } n_d} ]^T.$$

# Absorption Rate Measurements

- ▶ We also assume that  $S$  may be **approximated** by a function written as:

$$\tilde{S}(x,y) = \sum_{b=1}^B \alpha_b \tilde{s}_b(x,y);$$

- ▶ Where  $\tilde{s}_b(x,y)$ ,  $b = 1, \dots, B$ , corresponds to a set of basis functions defined in  $D$  and  $\alpha_b$  are the targets of the estimation process;
- ▶ If we try to directly relate  $r$  with the source  $\tilde{S}$ , we find out that we end up with a **nonlinear relation** between them due to the requirement of first solving the transport equation.

# Absorption Rate Measurements

- ▶ Alternatively, we can use the **adjoint of the transport equation** to write a new expression for the rate of absorption, which leads us to a **linear model** that relates  $r$  and  $\tilde{S}$ . Thus, we write [2]:

$$r = \langle \psi^\dagger, \tilde{S} \rangle - P[\psi, \psi^\dagger];$$

- ▶ Where  $\psi^\dagger$  is the **adjoint flux** and satisfy the **adjoint transport equation**:

$$\mathcal{L}^\dagger \psi^\dagger = \sigma_d;$$

- ▶ With  $\mathcal{L}^\dagger$  being the **adjoint transport operator** defined by:

$$\mathcal{L}^\dagger \psi^\dagger(\mathbf{r}, \mathbf{\Omega}) = -\mathbf{\Omega} \cdot \nabla \psi^\dagger + \sigma \psi^\dagger - \int_{\mathbb{S}'} \sigma_s \psi^\dagger d\mathbf{\Omega}'.$$

# Absorption Rate Measurements

- ▶ The  $P[\psi, \psi^\dagger]$  term depends on the **boundary conditions** of both transport problems;
- ▶ In fact, if we define boundary conditions at the **outward** directions of the adjoint problem:

$$\psi^\dagger(\cdot, \mathbf{\Omega}_{out}) = \rho\psi^\dagger(\cdot, \mathbf{\Omega}_{in});$$

- ▶ Then, the  $P[\psi, \psi^\dagger]$  term will depend only on the boundary conditions of the **original transport problem**, that is:

$$P[\psi, \psi^\dagger] = \int_{\partial V} \int_{\mathbf{\Omega} \cdot \mathbf{n} < 0} \mathbf{n} \cdot \mathbf{\Omega} \psi_b(\mathbf{r}, \mathbf{\Omega}) \psi^\dagger(\mathbf{r}, \mathbf{\Omega}) d\mu.$$

# Absorption Rate Measurements

- ▶ The absorption rate measurements vector is rewritten as:

$$\mathbf{r}(\boldsymbol{\alpha}) = \mathbf{A}\boldsymbol{\alpha} + \mathbf{p};$$

- ▶ Where  $\mathbf{A} = [A_{i,b}]$  is a  $D \times B$  matrix whose components we write as:

$$A_{b,i} = \langle \psi_i^\dagger, \tilde{s}_b \rangle;$$

- ▶  $\mathbf{p} = [p_i]$  is a  $D$ -dimensional vector whose components are  $p_i = P[\psi, \psi_i^\dagger]$ ;
- ▶ And  $\boldsymbol{\alpha} = [\alpha_b]$  is a  $B$ -dimensional vector with the targets of the estimation process;
- ▶ **Question:** how to solve the adjoint transport equation?

# ADO-Nodal Formulation

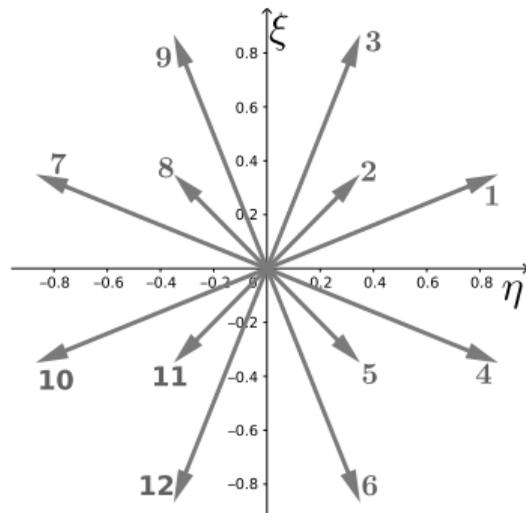
## 2D Adjoint Monoenergetic Problems

## ADO-Nodal Formulation to the Adjoint Transport Equation

- ▶ For each region, we consider the **averaged flux** with respect to the  $y$  variable [2]:

$$\psi_y^\dagger(x, \mathbf{\Omega}) = \frac{1}{h_y} \int_{y_{p+1}}^{y_p} \psi^\dagger(x, y, \mathbf{\Omega}) dy;$$

- ▶ We use  $M$  nodes,  $\mathbf{\Omega}_k = (\eta_k, \xi_k)$ , and weights,  $w_k$  of the projection of the known **LQ<sub>n</sub> quadrature** in the plane  $\eta O \xi$  [7];
- ▶ We consider discrete directions **ordered** according to:



## ADO-Nodal Formulation to the Adjoint Transport Equation

- ▶ **Integrating** the adjoint transport equation with respect to the  $y$  variable, we obtain:

$$\mp \eta_m \frac{d}{dx} \psi_y^\dagger(x, \Omega_m) + \sigma \psi_y^\dagger(x, \Omega_m) - \sigma_s \sum_{n=1}^{M/2} w_n \left[ \psi_y^\dagger(x, \Omega_n) + \psi_y^\dagger(x, \Omega_{n+M/2}) \right] = S_y^\dagger(x, \Omega_m);$$

- ▶ Where  $S_y^\dagger(x, \Omega_m)$  is such that:

$$S_y^\dagger(x, \Omega_m) = \frac{\xi_m}{h_y} \left[ \psi^\dagger(x, y_p, \Omega_m) - \psi^\dagger(x, y_{p+1}, \Omega_m) \right] + \frac{1}{h_y} \int_{y_{p+1}}^{y_p} S_y^\dagger(x, y, \Omega_m) dy;$$

- ▶ We assume that the **transversal leakage** terms  $\psi^\dagger(x, y_p, \Omega_m)$  and  $\psi^\dagger(x, y_{p+1}, \Omega_m)$  are constant for each direction.

## ADO-Nodal Formulation to the Adjoint Transport Equation

- ▶ We apply the known one-dimensional adjoint ADO methodology to obtain **homogeneous solutions** [8]:

$$\psi_{y,h}^\dagger(x, \mathbf{\Omega}_m) = \sum_{j=1}^{M/2} \left\{ A_j \phi_y(\nu_j, \mathbf{\Omega}_m) e^{-(x-x_q)/\nu_j} + A_{j+M/2} \phi_y(\nu_j, \mathbf{\Omega}_{m+M/2}) e^{-(x_{q+1}-x)/\nu_j} \right\},$$

$$\psi_{y,h}^\dagger(x, \mathbf{\Omega}_{m+M/2}) = \sum_{j=1}^{M/2} \left\{ A_j \phi_y(\nu_j, \mathbf{\Omega}_{m+M/2}) e^{-(x-x_q)/\nu_j} + A_{j+M/2} \phi_y(\nu_j, \mathbf{\Omega}_m) e^{-(x_{q+1}-x)/\nu_j} \right\};$$

- ▶ **Constant source** for each region  $\Rightarrow$  **constant particular solution** for each region:

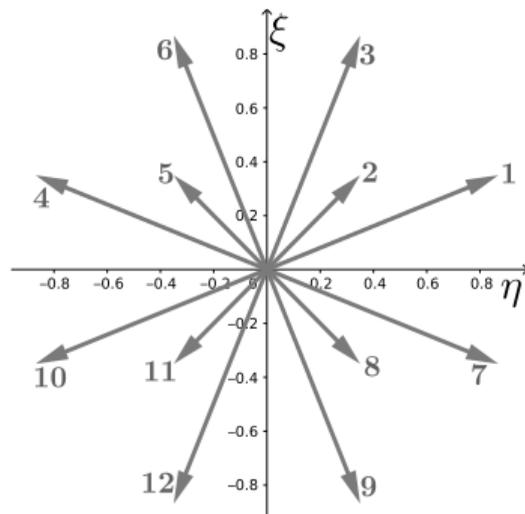
$$\sigma K_m - \sigma_s \sum_{n=1}^M w_n K_n = S_y^\dagger + \frac{\xi_m}{h_y} \left( \psi^\dagger(x, y_p, \mathbf{\Omega}_m) - \psi^\dagger(x, y_{p+1}, \mathbf{\Omega}_m) \right).$$

## ADO-Nodal Formulation to the Adjoint Transport Equation

- ▶ Similarly, we consider the **averaged flux** with respect to the  $x$  variable:

$$\psi_x^\dagger(y, \mathbf{\Omega}) = \frac{1}{h_x} \int_{x_q}^{x_{q+1}} \psi^\dagger(x, y, \mathbf{\Omega}) dx;$$

- ▶ We still use  $M$  nodes,  $\mathbf{\Omega}_k = (\eta_k, \xi_k)$ , and weights,  $w_k$  of the projection of the known **LQ<sub>n</sub> quadrature** in the plane  $\eta O \xi$ ;
- ▶ But, **this time**, we consider discrete directions **ordered** according to:



## ADO-Nodal Formulation to the Adjoint Transport Equation

- ▶ This way, **integrating** the transport equation with respect to  $\mathbf{x}$ , we obtain:

$$\mp \xi_m \frac{d}{dy} \psi_x^\dagger(y, \Omega_m) + \sigma \psi_x^\dagger(y, \Omega_m) - \sigma_s \sum_{n=1}^{M/2} w_n \left[ \psi_x^\dagger(y, \Omega_n) + \psi_x^\dagger(y, \Omega_{n+M/2}) \right] = S_x^\dagger(y, \Omega_m);$$

- ▶ Where  $S_x^\dagger(y, \Omega_m)$  is such that:

$$S_x^\dagger(y, \Omega_m) = \frac{\eta_m}{h_x} \left[ \psi^\dagger(x_{q+1}, y, \Omega_m) - \psi^\dagger(x_q, y, \Omega_m) \right] + \frac{1}{h_x} \int_{x_q}^{x_{q+1}} S_x^\dagger(x, y, \Omega_m) dx;$$

- ▶ We also assume constant **transversal leakage** terms  $\psi^\dagger(x_q, y, \Omega_m)$  and  $\psi^\dagger(x_{q+1}, y, \Omega_m)$

## ADO-Nodal Formulation to the Adjoint Transport Equation

- ▶ Finally, we apply the one-dimensional ADO methodology to obtain **homogeneous solutions**:

$$\psi_{x,h}^\dagger(y, \mathbf{\Omega}_m) = \sum_{j=1}^{M/2} \{ B_j \phi_x(\gamma_j, \mathbf{\Omega}_m) e^{-(y-y_{p+1})/\gamma_j} + B_{j+M/2} \phi_x(\gamma_j, \mathbf{\Omega}_{m+M/2}) e^{-(y_p-y)/\gamma_j} \},$$

$$\psi_{x,h}^\dagger(y, \mathbf{\Omega}_{m+M/2}) = \sum_{j=1}^{M/2} \{ B_j \phi_x(\gamma_j, \mathbf{\Omega}_{m+M/2}) e^{-(y-y_{p+1})/\gamma_j} + B_{j+M/2} \phi_x(\gamma_j, \mathbf{\Omega}_m) e^{-(y_p-y)/\gamma_j} \};$$

- ▶ **Constant source** for each region  $\Rightarrow$  **constant particular solution** for each region:

$$\sigma W_m - \sigma_s \sum_{n=1}^M w_n W_n = S_x^\dagger + \frac{\eta_m}{h_x} \left( \psi^\dagger(x_{q+1}, y, \mathbf{\Omega}_m) - \psi^\dagger(x_q, y, \mathbf{\Omega}_m) \right).$$

# ADO-Nodal Formulation to the Adjoint Transport Equation

- ▶ To obtain the coefficients  $A_j$  and  $B_j$  we deduce a **linear system** obtained by considering:
  - (i) **Particular solutions** for the averaged fluxes within each region:
    - (a) Regions that do not contain the boundaries – **constant leakage terms**;
    - (b) Regions that contain the boundaries – **boundary conditions**.
  - (ii) **Boundary conditions** for the averaged fluxes within each region;
  - (iii) **Continuity** of the averaged fluxes at the interfaces of two **contiguous regions**.
- ▶ We obtain a system of order  $4Mn_xn_y$ .

# Explicit Formulas for the Absorption Rate Measurements

# ADO-Nodal Formulation to the Adjoint Transport Equation

## Explicit Formulas for the Absorption Rate

- ▶ As previously stated, the **rate of absorption** is given by:

$$r = \langle \psi^\dagger, \tilde{S} \rangle - P[\psi, \psi^\dagger];$$

- ▶ Where we may use the **explicit expressions for the averaged fluxes** to write:

$$\langle \psi^\dagger, \tilde{S} \rangle = \sum_{q=1}^{n_x} \sum_{p=1}^{n_y} \tilde{S}_{p,q} h_{x,p,q} h_{y,p,q} \phi_{p,q}^\dagger.$$

# ADO-Nodal Formulation to the Adjoint Transport Equation

## Explicit Formulas for the Absorption Rate

- ▶  $\phi_{p,q}^\dagger$  represents the **averaged scalar flux** of particles, and is given by either:

$$\phi_{p,q}^\dagger = \frac{1}{h_{x,p,q}} \sum_{j=1}^{M/2} \nu_j^{p,q} \left( 1 - e^{-(x_{q+1}-x_q)/\nu_j^{p,q}} \right) \times \\ \left( A_j^{p,q} + A_{j+M/2}^{p,q} \right) \phi_{y,j}^{0,p,q} + \phi_{y,j}^{1,p,q};$$

- ▶ Or:

$$\phi_{p,q}^\dagger = \frac{1}{h_{y,p,q}} \sum_{j=1}^{M/2} \gamma_j^{p,q} \left( 1 - e^{-(y_p-y_{p+1})/\gamma_j^{p,q}} \right) \times \\ \left( B_j^{p,q} + B_{j+M/2}^{p,q} \right) \phi_{x,j}^{0,p,q} + \phi_{x,j}^{1,p,q};$$

- ▶ Where  $\phi_{y,j}^{0,p,q}$ ,  $\phi_{y,j}^{1,p,q}$ ,  $\phi_{x,j}^{0,p,q}$  and  $\phi_{x,j}^{1,p,q}$  are closed form terms and depend on the eigenfunctions and internal sources.

# ADO-Nodal Formulation to the Adjoint Transport Equation

## Explicit Formulas for the Absorption Rate

- ▶ Further, We write the boundary term  $P[\psi, \psi^\dagger]$  as:

$$P[\psi, \psi^\dagger] = P_1[\psi, \psi^\dagger] + P_2[\psi, \psi^\dagger] + P_3[\psi, \psi^\dagger] + P_4[\psi, \psi^\dagger];$$

- ▶ Where:

$$P_1[\psi, \psi^\dagger] = - \sum_{q=1}^{n_x} \sum_{m=1}^{M/2} w_m h_{x, n_y, q} f_{n_y, q} \psi_{x, n_y, q}^\dagger(0, \Omega_m),$$

$$P_2[\psi, \psi^\dagger] = - \sum_{q=1}^{n_x} \sum_{m=1}^{M/2} w_m h_{x, 1, q} f_{1, q} \psi_{x, 1, q}^\dagger(b, \Omega_{m+M/2}),$$

$$P_3[\psi, \psi^\dagger] = - \sum_{p=1}^{n_y} \sum_{m=1}^{M/2} w_m h_{y, 1, q} f_{p, 1} \psi_{y, p, 1}^\dagger(0, \Omega_m),$$

$$P_4[\psi, \psi^\dagger] = - \sum_{p=1}^{n_y} \sum_{m=1}^{M/2} w_m h_{y, p, n_x} f_{p, n_x} \psi_{y, p, n_x}^\dagger(a, \Omega_{m+M/2}).$$

# Source Estimation

## Source Estimation

- ▶ Given **exact** measurements  $\mathbf{r}(\boldsymbol{\alpha})$ , we consider **additive noise** such that:

$$\tilde{\mathbf{r}}_m = \mathbf{r}(\boldsymbol{\alpha}) + \boldsymbol{\epsilon};$$

- ▶ If  $\boldsymbol{\epsilon} \sim \mathcal{N}(\mathbf{0}, \mathbf{W})$ , with  $\mathbf{W} = \sigma^2 \mathbf{I}$ , it is possible to write the PDF for  $\boldsymbol{\epsilon}$  as [6]:

$$\pi(\tilde{\mathbf{r}}|\boldsymbol{\alpha}) = (2\pi)^{-n_d/2} |\mathbf{W}|^{-1/2} \exp \left\{ -\frac{1}{2} [\tilde{\mathbf{r}} - \mathbf{r}]^T \mathbf{W}^{-1} [\tilde{\mathbf{r}} - \mathbf{r}] \right\};$$

- ▶ Which is **maximized** when the maximum likelihood function

$$S_{ML}(\boldsymbol{\alpha}) = [\tilde{\mathbf{r}} - \mathbf{r}]^T [\tilde{\mathbf{r}} - \mathbf{r}] = \|\mathbf{A}\boldsymbol{\alpha} - \mathbf{b}\|^2,$$

is **minimized**.

## Source Estimation

- ▶ We obtain **exact results** when using **noise free** measurements calculated using the ADO method with the **same valor of  $N$**  and the **same mesh** used to calculate  $\mathbf{A}$ ;
- ▶ However, **as expected**, we got **oscillations** on our estimates when considering **noisy measurements**;
- ▶ **Tikhonov regularization**: the maximum likelihood function is modified according to:

$$S_{\lambda}(\boldsymbol{\alpha}) = \|\mathbf{A}\boldsymbol{\alpha} - \mathbf{b}\|^2 + \lambda \|\boldsymbol{\alpha} - \boldsymbol{\alpha}_0\|^2;$$

- ▶ **Iterated Tikhonov** to find the best values of  $\lambda$  (discrepancy principle).

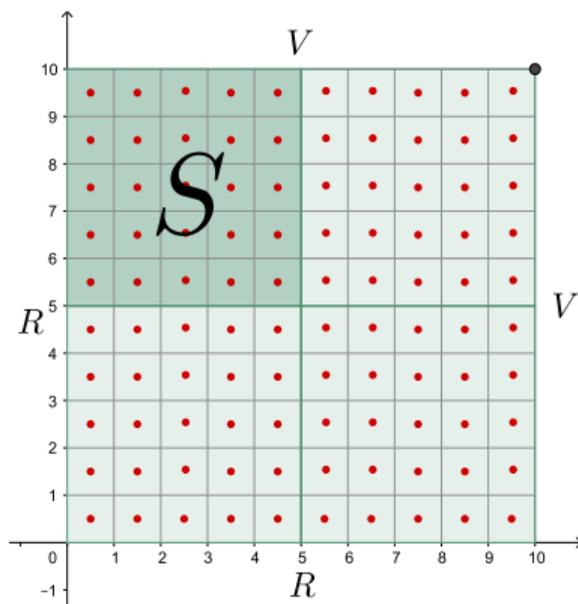
# Source Estimation

## Test Problem I

# Source Estimation – Test Problem I

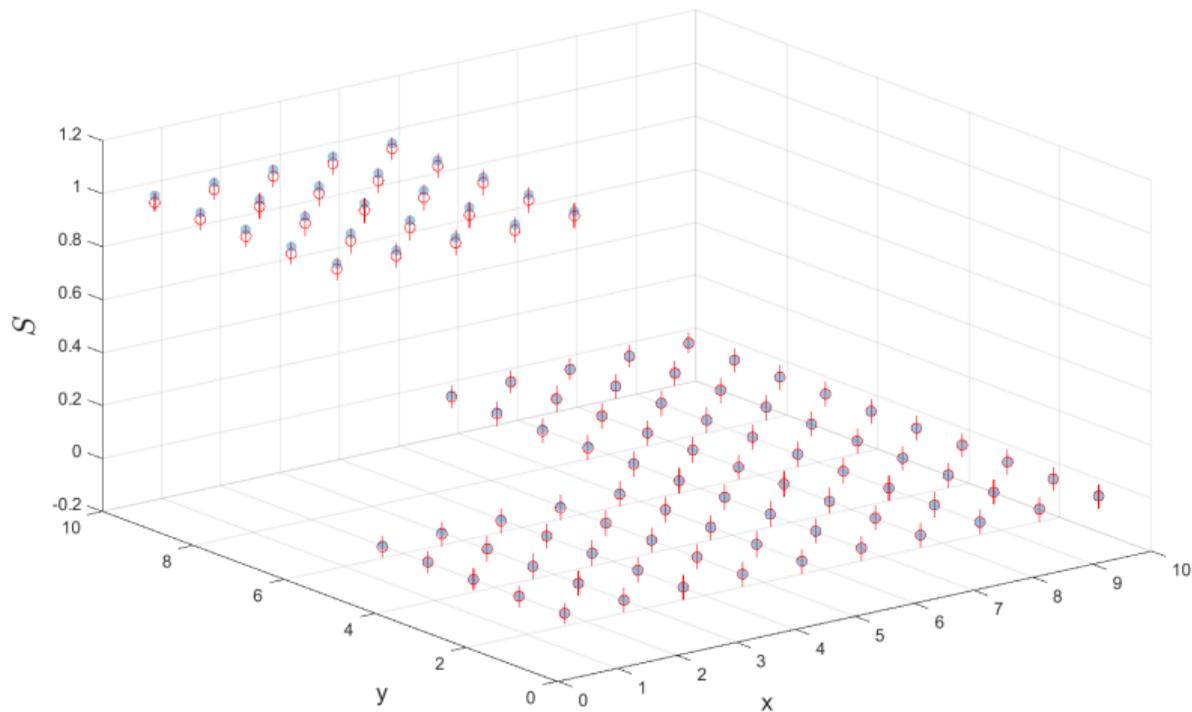
## Test Problem I configuration

- ▶ Rectangle  $(x, y) \in [0, 10] \times [0, 10]$ ;
- ▶ Heterogeneous medium:
$$\sigma_t(x, y) = \begin{cases} 1.0, & \text{for } (x, y) \in [0, 5] \times [0, 5], \\ 2.0, & \text{otherwise;} \end{cases}$$
$$\sigma_s(x, y) = \begin{cases} 0.5, & \text{for } (x, y) \in [0, 5] \times [0, 5], \\ 0.1, & \text{otherwise;} \end{cases}$$
- ▶ Internal source of particles:
$$S(x, y) = \begin{cases} 1.0, & \text{for } (x, y) \in [0, 5] \times [6, 10], \\ 0.0, & \text{otherwise;} \end{cases}$$
- ▶ Diamond difference to generate “exact” measurements, then add 1% of additive white noise.



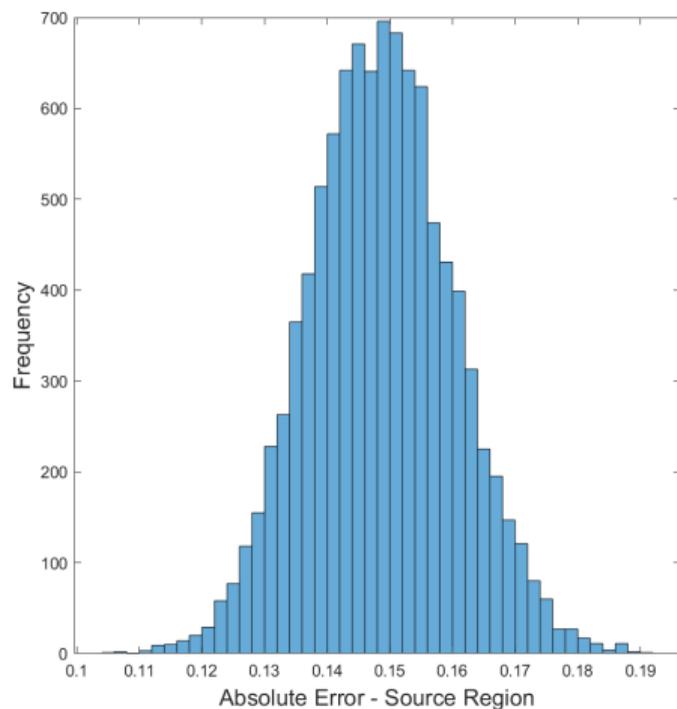
# Source Estimation – Test Problem I

Mean value of estimates and confidence intervals



# Source Estimation – Test Problem I

## Errors



- ▶ Minimum, mean, maximum and standard deviation of the absolute errors ( $L^2$  norm):

Errors	Min.	Mean	Max.	Std.
$\ \tilde{S} - \hat{S}\ $	0.1054	0.1487	0.1914	0.0116

- ▶ 99% confidence interval:  
(0.1193, 0.1796)

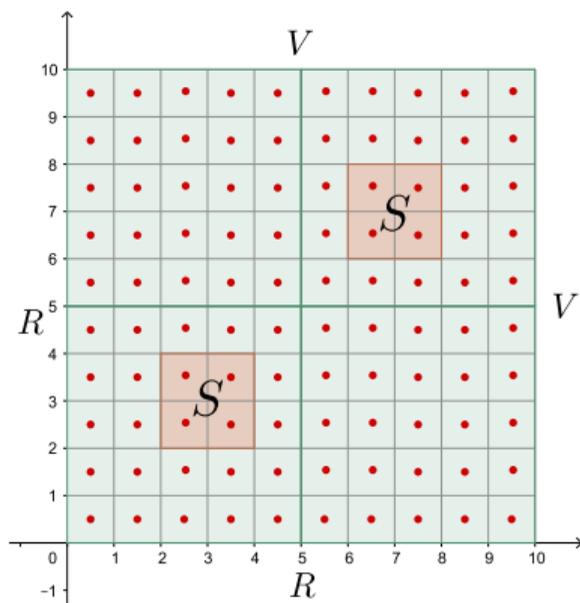
# Source Estimation

## Test Problem II

# Source Estimation – Test Problem II

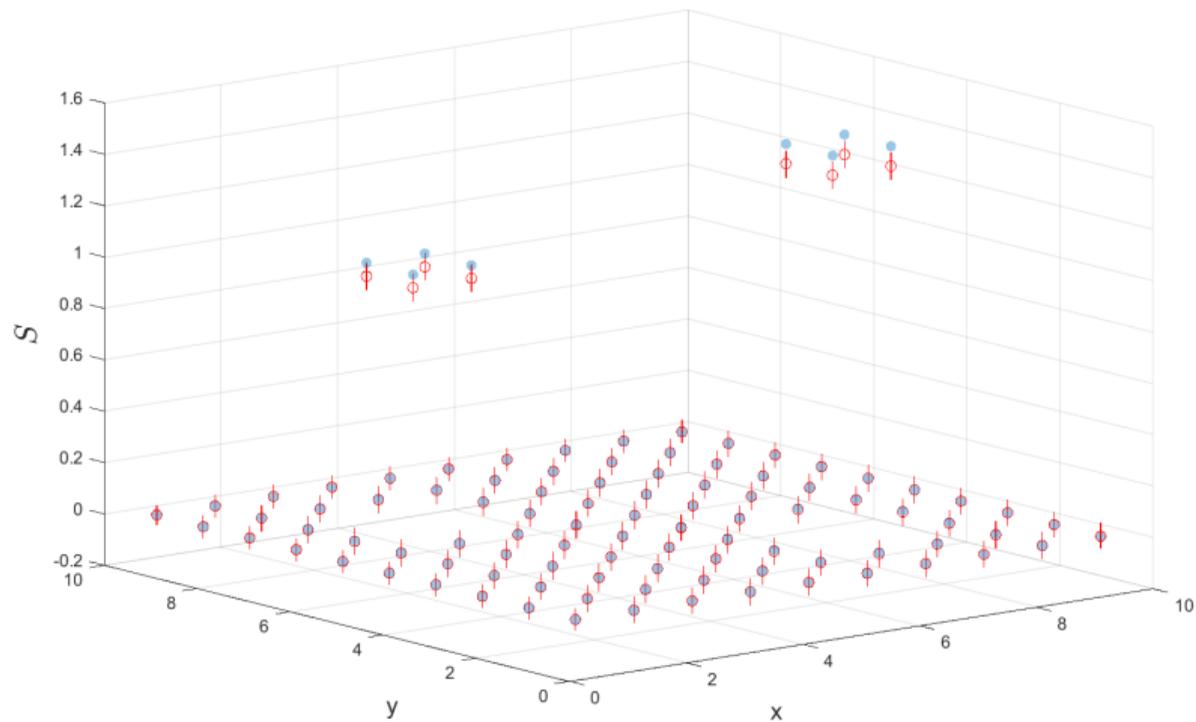
## Test Problem II configuration

- ▶ Rectangle  $(x, y) \in [0, 10] \times [0, 10]$ ;
- ▶ Heterogeneous medium:
$$\sigma_t(x, y) = \begin{cases} 1.0, & \text{for } (x, y) \in [0, 5] \times [0, 5], \\ 2.0, & \text{otherwise;} \end{cases}$$
$$\sigma_s(x, y) = \begin{cases} 0.5, & \text{for } (x, y) \in [0, 5] \times [0, 5], \\ 0.1, & \text{otherwise;} \end{cases}$$
- ▶ Internal source of particles:
$$S(x, y) = \begin{cases} 1.0, & \text{for } (x, y) \in [2, 4] \times [2, 4]; \\ 1.5, & \text{for } (x, y) \in [6, 8] \times [6, 8]; \\ 0.0, & \text{otherwise;} \end{cases}$$
- ▶ Diamond difference to generate “exact” measurements, then add 1% of additive white noise.



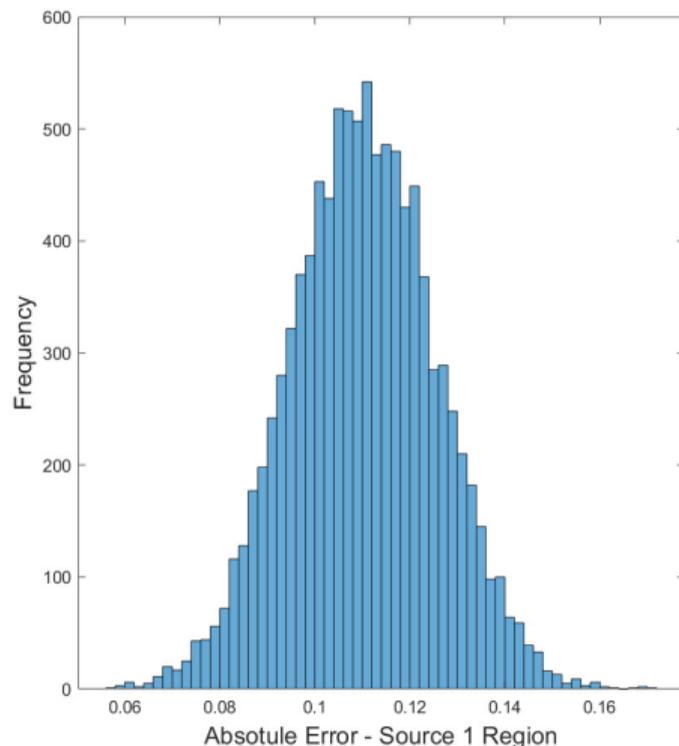
# Source Estimation – Test Problem II

Mean value of estimates and confidence intervals



# Source Estimation – Test Problem II

## Errors – $S_1$ region



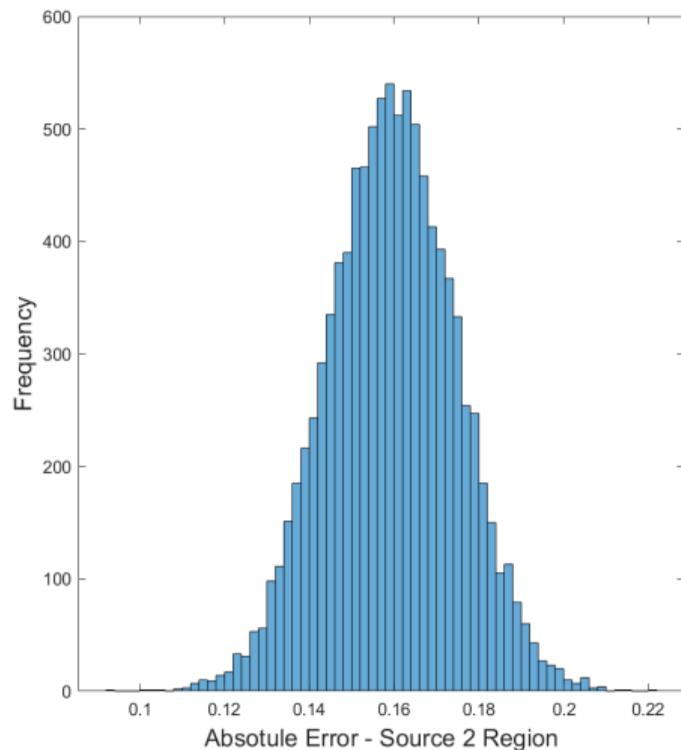
- ▶ Minimum, mean, maximum and standard deviation of the absolute errors ( $L^2$  norm) at  $S_1 = 1.0$  region:

Errors	Min.	Mean	Max.	Std.
$\ \tilde{S}_1 - \hat{S}_1\ $	0.0569	0.1102	0.1717	0.0153

- ▶ 99% confidence interval:  
(0.0701, 0.1488)

# Source Estimation – Test Problem II

## Errors – $S_2$ region



- ▶ Minimum, mean, maximum and standard deviation of the absolute errors ( $L^2$  norm) at  $S_2 = 1.5$  region:

Errors	Min.	Mean	Max.	Std.
$\ \tilde{S}_2 - \hat{S}_2\ $	0.0920	0.1595	0.2208	0.0151

- ▶ 99% confidence interval:  
(0.1200, 0.1988)

# Bayesian Inference Techniques

# Bayesian Inference Techniques

- ▶ Using **Bayes' theorem** in  $\pi_{post}(\boldsymbol{\alpha}) = \pi(\boldsymbol{\alpha}|\tilde{\mathbf{r}})$ :

$$\pi_{post}(\boldsymbol{\alpha}) = \frac{\pi(\tilde{\mathbf{r}}|\boldsymbol{\alpha})\pi(\boldsymbol{\alpha})}{\pi(\tilde{\mathbf{r}})} \propto \pi(\tilde{\mathbf{r}}|\boldsymbol{\alpha})\pi(\boldsymbol{\alpha});$$

- ▶ Prior distribution  $\pi(\boldsymbol{\alpha})$  for the parameters – **Gaussian Markov random field**;
- ▶ Parameter  $\lambda$  treated as a hyperparameter – **Rayleigh's distribution**;
- ▶ Posterior distribution of probability:

$$\pi_{post}(\boldsymbol{\alpha}, \lambda) \propto \lambda^{\frac{N_d+2}{2}} \exp \left\{ -\frac{1}{2} [\tilde{\mathbf{r}} - \mathbf{r}(\boldsymbol{\alpha})]^T \mathbf{W}^{-1} [\tilde{\mathbf{r}} - \mathbf{r}(\boldsymbol{\alpha})] - \frac{1}{2} \lambda [\boldsymbol{\alpha} - \tilde{\boldsymbol{\alpha}}]^T \mathbf{Z} [\boldsymbol{\alpha} - \tilde{\boldsymbol{\alpha}}] - \frac{1}{2} \left( \frac{\lambda}{\lambda_0} \right)^2 \right\}.$$

# Bayesian Inference Techniques

- ▶ We use the **Metropolis-Hastings algorithm** [6] to sample the posterior distribution of the source parameters:
  - ▶ Consider an **initial state**  $\alpha^{(0)}, \lambda^{(0)}$ ;
  - ▶ Given a state  $\alpha^{(k-1)}, \lambda^{(k-1)}$ , we perform a **random walk** step to find a **proposal state**  $\alpha^{(*)}, \lambda^{(*)}$  given by:

$$\alpha^{(*)} = \alpha^{(k-1)} + \alpha, \quad \lambda^{(*)} = \lambda^{(k-1)} + \lambda;$$

- ▶ We **accept** the new state, that is, we set  $\alpha^{(k)} = \alpha^{(*)}$  and  $\lambda^{(k)} = \lambda^{(*)}$ , if [6]:

$$\min \left\{ 1, \frac{\pi_{post}(\alpha^{(*)}, \lambda^{(*)})}{\pi_{post}(\alpha^{(k-1)}, \lambda^{(k-1)})} \right\} \geq \mu,$$

where  $\mu \sim \mathcal{U}(0, 1)$ ;

- ▶ Otherwise, we **reject** the new state, that is, we set  $\alpha^{(k)} = \alpha^{(k-1)}$  and  $\lambda^{(k)} = \lambda^{(k-1)}$ .

# Source Estimation

## Test Problem III

# Source Estimation – Test Problem III

## Test Problem III configuration

- ▶ Same physical properties of the Test Problem I;
- ▶ Rectangle  $(x, y) \in [0, 10] \times [0, 10]$ ;

- ▶ Heterogeneous medium:

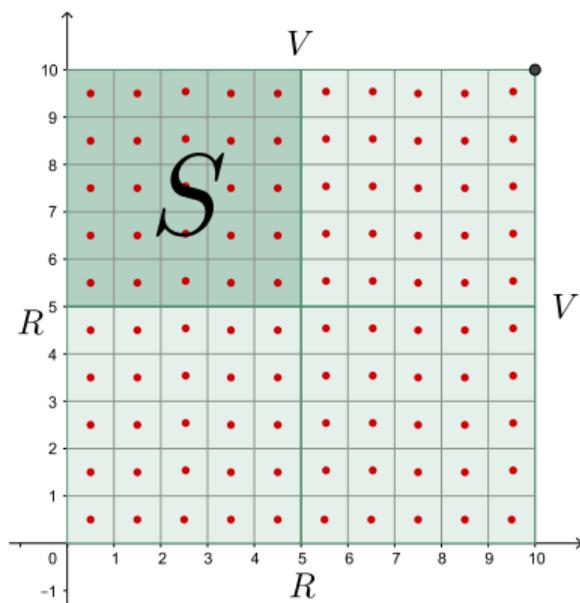
$$\sigma_t(x, y) = \begin{cases} 1.0, & \text{for } (x, y) \in [0, 5] \times [0, 5], \\ 2.0, & \text{otherwise;} \end{cases}$$

$$\sigma_s(x, y) = \begin{cases} 0.5, & \text{for } (x, y) \in [0, 5] \times [0, 5], \\ 0.1, & \text{otherwise;} \end{cases}$$

- ▶ Internal source of particles:

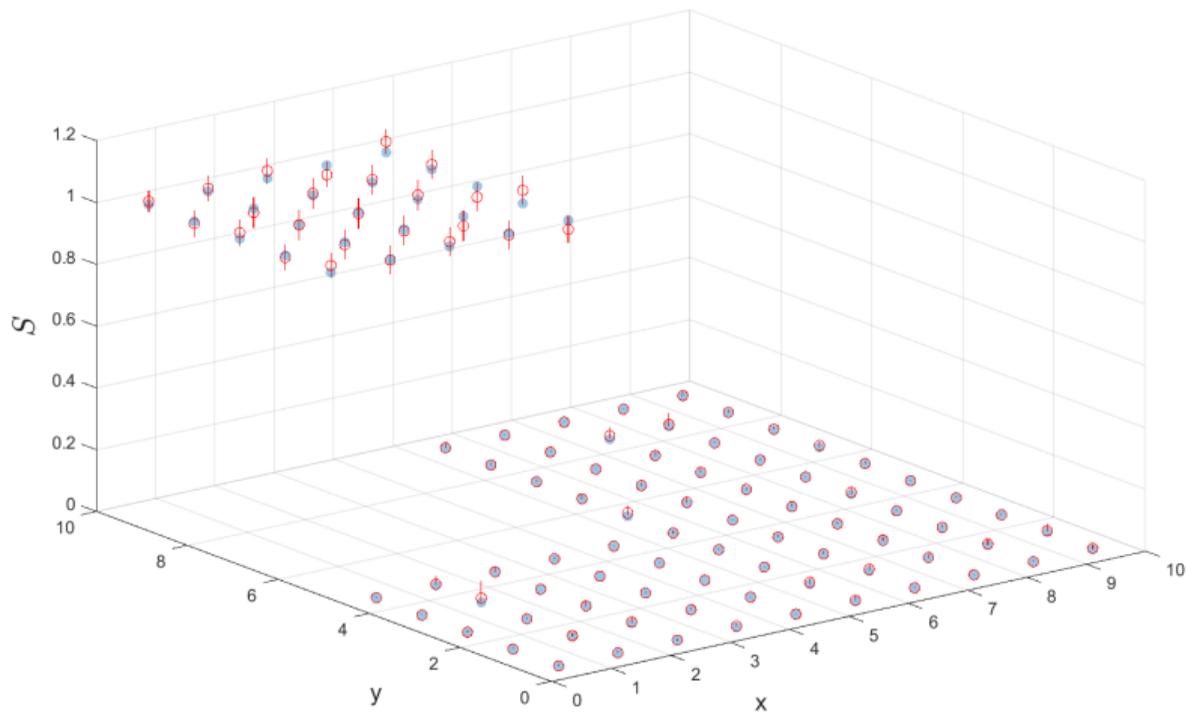
$$S(x, y) = \begin{cases} 1.0, & \text{for } (x, y) \in [0, 5] \times [6, 10], \\ 0.0, & \text{otherwise;} \end{cases}$$

- ▶ Diamond difference to generate “exact” measurements, then add 1% of additive white noise.



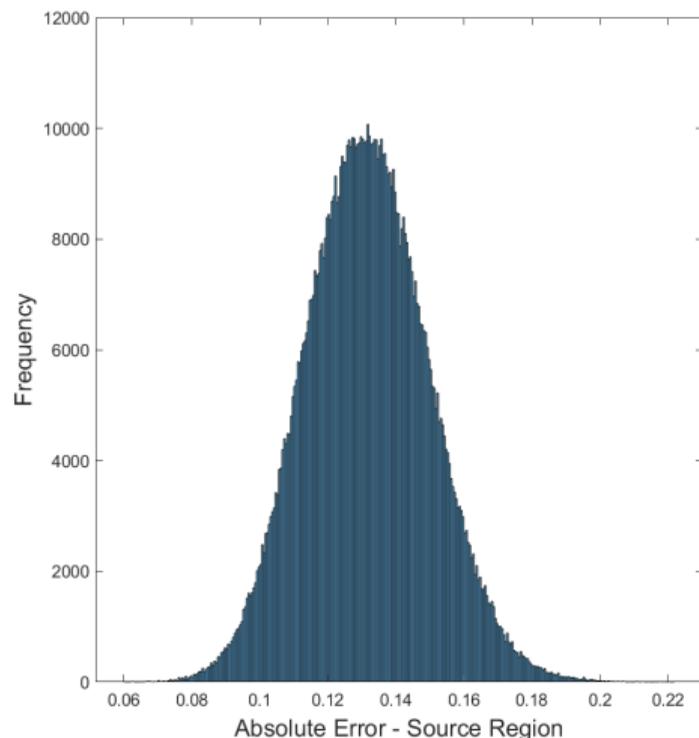
# Source Estimation – Test Problem III

Mean value of estimates and credible intervals



# Source Estimation – Test Problem III

## Errors



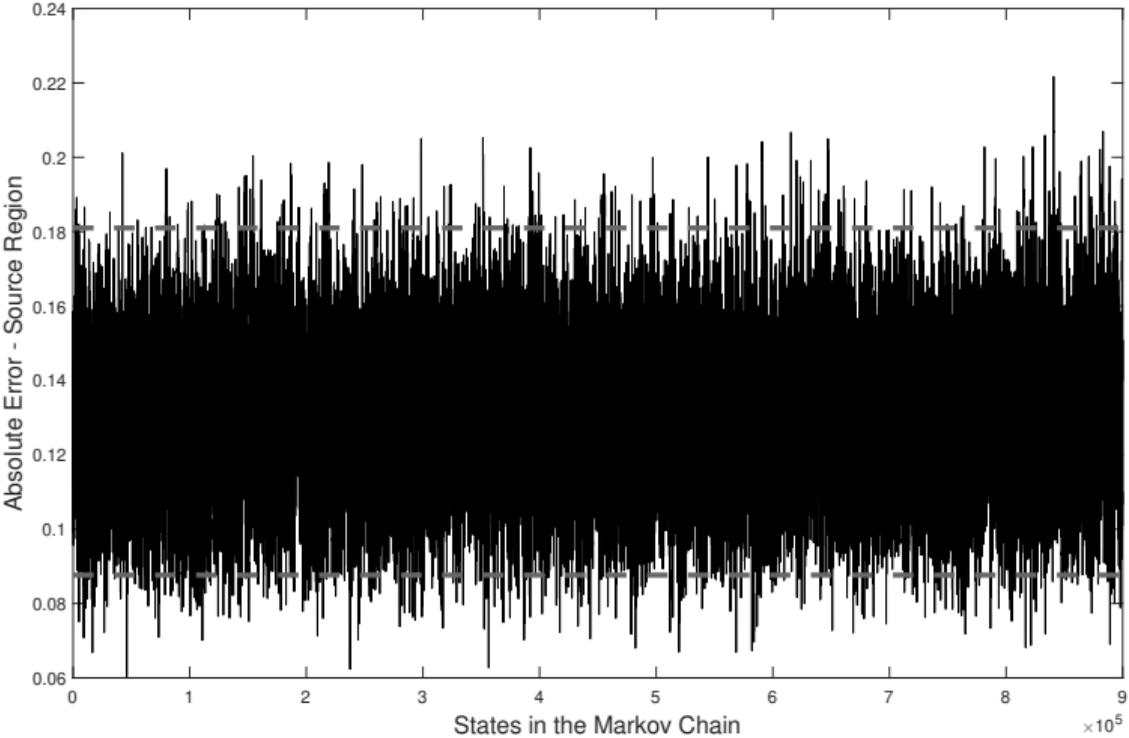
- ▶ Minimum, mean, maximum and standard deviation of the absolute errors ( $L^2$  norm):

Errors	Min.	Mean	Max.	Std.
$\ \tilde{S} - \hat{S}\ $	0.0603	0.1316	0.2217	0.0181

- ▶ 99% credible interval:  
(0.0876, 0.1810)

# Source Estimation – Test Problem III

States in the Markov chain and 99% credible interval



# Source Estimation

## Test Problem IV

# Source Estimation – Test Problem IV

## Test Problem IV configuration

- ▶ Same physical properties of the Test Problem II;

- ▶ Rectangle  $(x, y) \in [0, 10] \times [0, 10]$ ;

- ▶ Heterogeneous medium:

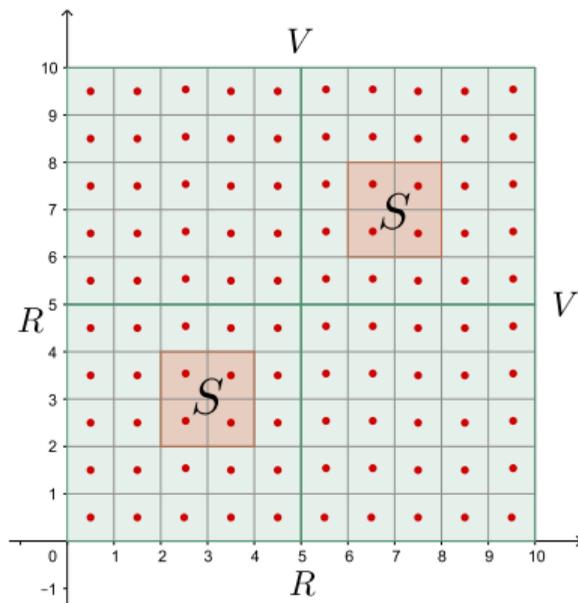
$$\sigma_t(x, y) = \begin{cases} 1.0, & \text{for } (x, y) \in [0, 5] \times [0, 5], \\ 2.0, & \text{otherwise;} \end{cases}$$

$$\sigma_s(x, y) = \begin{cases} 0.5, & \text{for } (x, y) \in [0, 5] \times [0, 5], \\ 0.1, & \text{otherwise;} \end{cases}$$

- ▶ Internal source of particles:

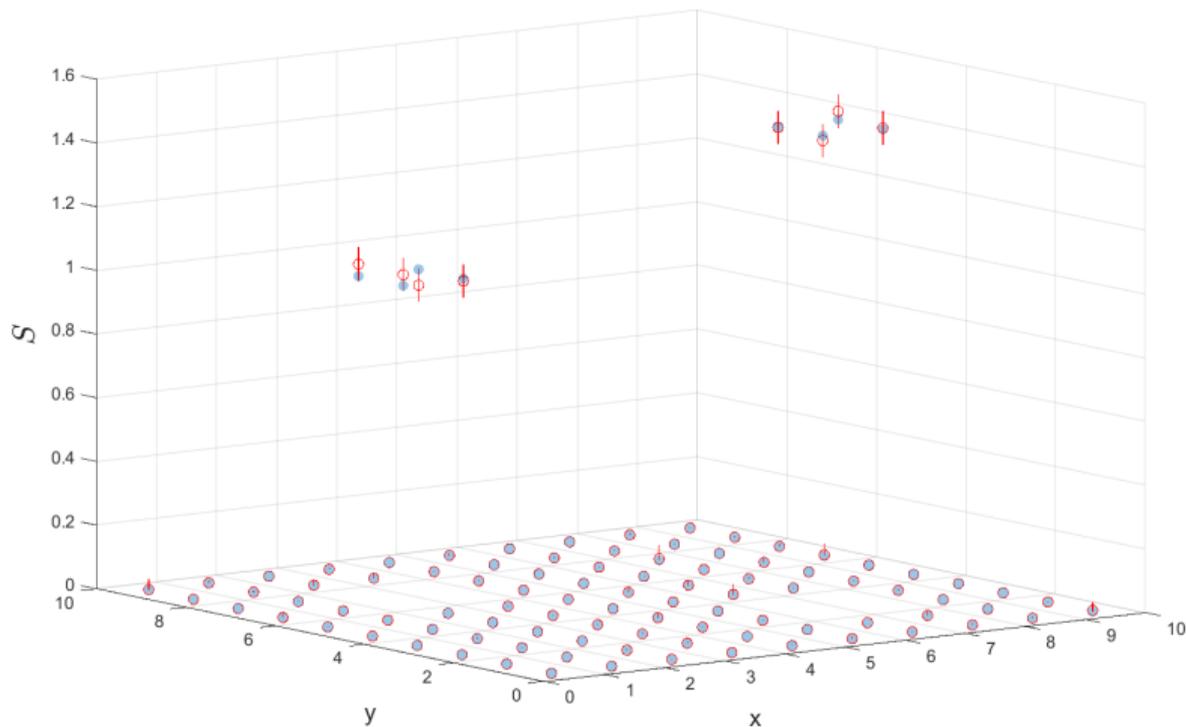
$$S(x, y) = \begin{cases} 1.0, & \text{for } (x, y) \in [2, 4] \times [2, 4]; \\ 1.5, & \text{for } (x, y) \in [6, 8] \times [6, 8]; \\ 0.0, & \text{otherwise;} \end{cases}$$

- ▶ Diamond difference to generate “exact” measurements, then add 1% of additive white noise.



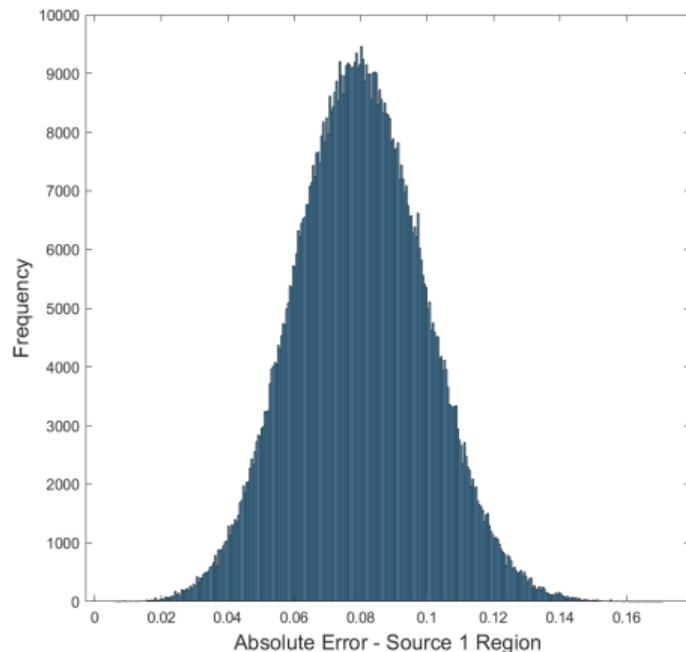
# Source Estimation – Test Problem IV

Mean value of estimates and credible intervals



# Source Estimation – Test Problem IV

## Errors – $S_1$ region



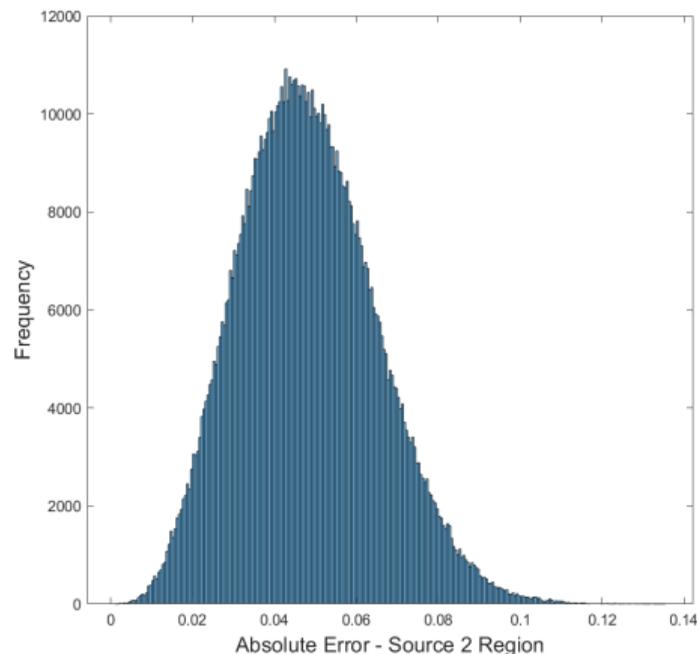
- ▶ Minimum, mean, maximum and standard deviation of the absolute errors ( $L^2$  norm) at  $S_1 = 1.0$  region:

Errors	Min.	Mean	Max.	Std.
$\ \tilde{S} - \hat{S}\ $	0.0056	0.0799	0.1708	0.0195

- ▶ 99% credible interval:  
(0.0319, 0.1313)

# Source Estimation – Test Problem IV

Errors –  $S_2$  region



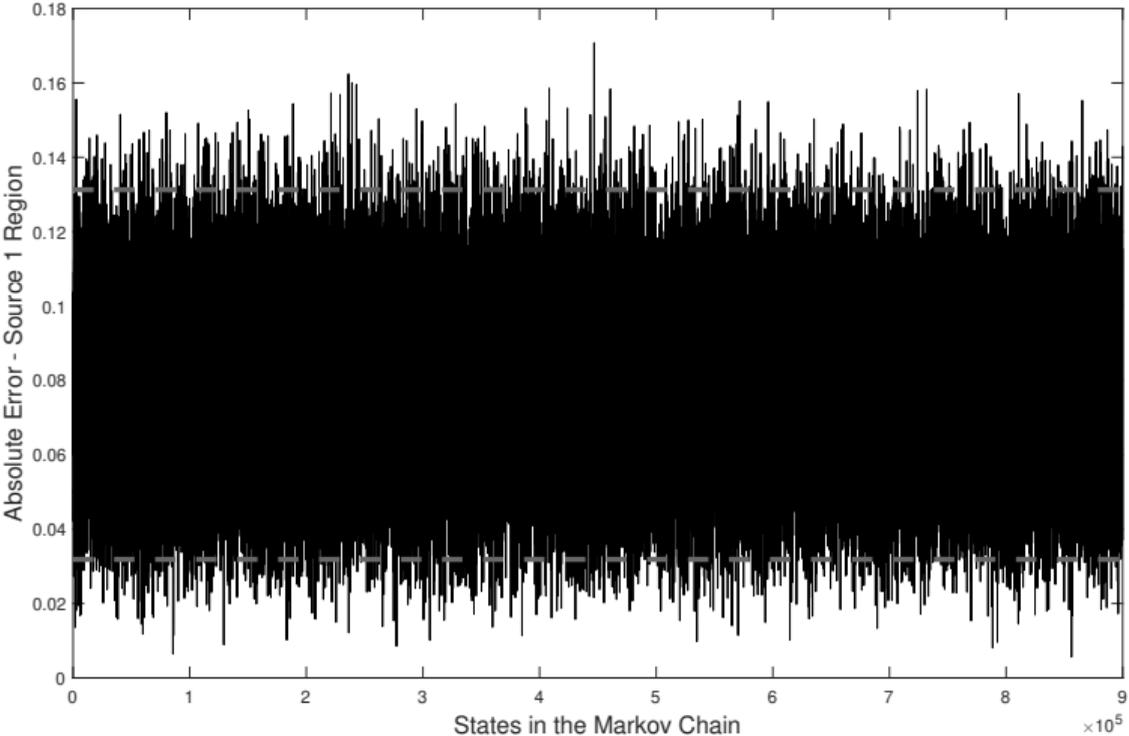
- ▶ Minimum, mean, maximum and standard deviation of the absolute errors ( $L^2$  norm) at  $S_2 = 1.5$  region:

Errors	Min.	Mean	Max.	Std.
$\ \tilde{S} - \hat{S}\ $	0.0013	0.0481	0.1353	0.0167

- ▶ 99% credible interval:  
(0.0122, 0.0955)

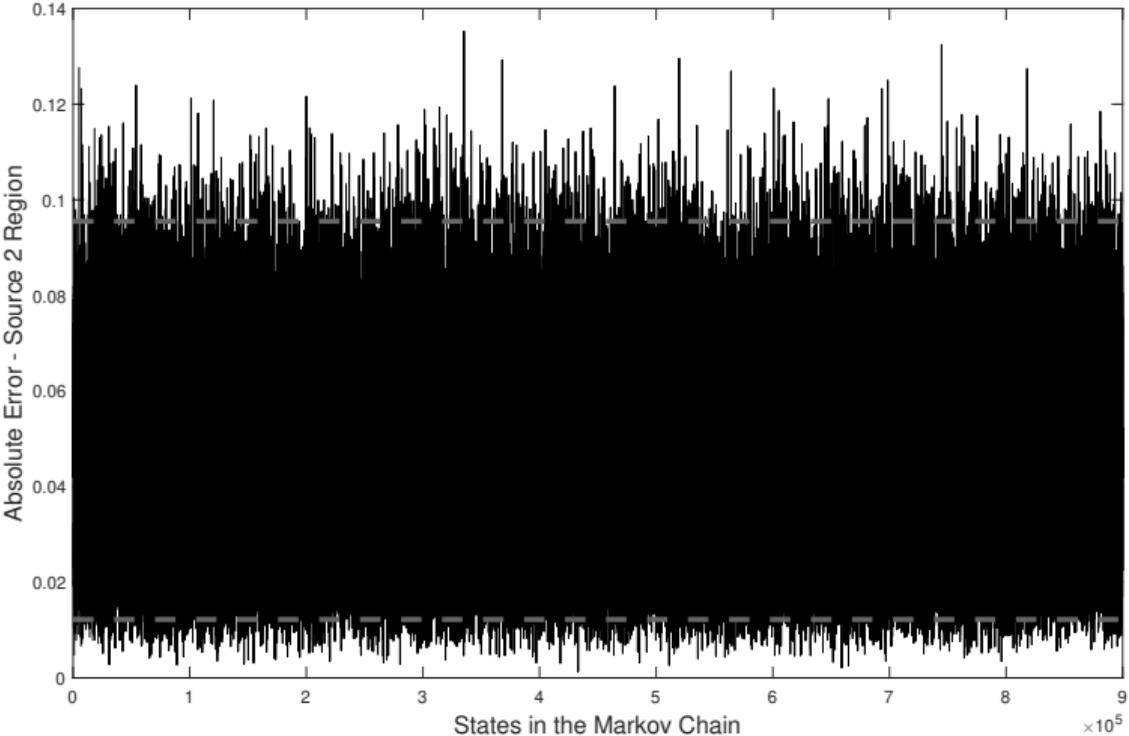
# Source Estimation – Test Problem IV

States in the Markov chain and 99% credible interval –  $S_1$  region



# Source Estimation – Test Problem IV

States in the Markov chain and 99% credible interval –  $S_2$  region



## Conclusions

- ▶ We proposed an **ADO-Nodal** solution for the adjoint transport equation in a rectangular region, considering the monoenergetic approximation and isotropic scattering;
- ▶ **Explicit solutions** allowed us to derive explicit expressions for  $r$ , which lead us to **speed and accuracy** when estimating the absorption rates;
- ▶ With respect to the **source estimation problems**, we believe to have attained good approximations in all cases;
- ▶ We found the results obtained considering the MCMC algorithm to be **slightly better** than the ones obtained using IT algorithm, in the sense of producing smaller absolute errors;
- ▶ Moreover, the estimated magnitudes **outside the sources region** were also better approximated using the Bayesian formulation;
- ▶ Finally, we found that the MCMC algorithm performed at least **twice as fast** as the IT, though no analysis has yet been performed.

# Publications

- ▶ C. B. Pazinato, S. R. Cromianski, R. C. Barros e L. B. Barichello, An Analytical Discrete Ordinates Solution for One-Speed Slab Geometry Adjoint Transport Problems with Isotropic Scattering. **M&C**, 2015.
- ▶ C. B. Pazinato, R. C. Barros e L. B. Barichello, *An Analytical Discrete Ordinates Formulation for Monoenergetic Slab-Geometry Source-Detector Calculations*. **International Journal of Nuclear Energy Science and Technology**, 2016.
- ▶ C. B. Pazinato e L. B. Barichello, *On the use of Analytical Techniques for Source Reconstruction Problems*. **M&C**, 2017.
- ▶ C. B. Pazinato e L. B. Barichello, *An Analytical Discrete Ordinates Solution for Multigroup Adjoint Transport Problems*. **ICTT**, 2017.
- ▶ C. B. Pazinato e L. B. Barichello, *On the Use of the Adjoint Operator for Source Reconstruction in Particle Transport Problems*. **Inverse Problems in Science and Engineering**, 2018.
- ▶ C. B. Pazinato e L. B. Barichello, *Energy Dependent Source Reconstructions via Explicit Formulations of the Adjoint Particles Flux*. **Journal of Computational and Theoretical Transport**, 2018.
- ▶ L. B. Barichello, K. Rui e C. B. Pazinato, *Advances in the ADO-Nodal Method for the Solution of Particle Transport Problems in Two Dimensional Media*. **PHYTRA**, 2018.
- ▶ L. B. Barichello, C. B. Pazinato e K. Rui, An analytical discrete ordinates nodal solution to the two-dimensional adjoint transport problem. **Annals of Nuclear Energy**, 2020.

## Ongoing works

- ▶ More detailed study of **Bayesian inference** techniques for solving source estimation inverse problems;
- ▶ Use of more refined meshes – analysis of the ADO-Nodal matrices;
- ▶ Energy dependence – multigroup approximation;
- ▶ Arbitrary degree of anisotropy.

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Thank you!

# On the use of the Adjoint Operator in Estimating Neutral Particles Sources

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